

OPTICAL ANALYZER AND METHOD FOR MEASURING SPECTRAL  
AMPLITUDE AND PHASE OF INPUT OPTICAL SIGNALS USING  
HETERODYNE ARCHITECTURE

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FIELD OF THE INVENTION

[0001] The invention relates generally to optical analyzers, and more particularly to an optical analyzer and method for measuring spectral amplitude  
10 and phase of input optical signals.

BACKGROUND OF THE INVENTION

[0002] Spectral amplitude and phase measurements of optical signals are often  
15 desired to characterize the signals in the time domain. These spectral measurements allow time varying optical characteristics of the optical signals to be studied by means of the Fourier transform. An example of a time varying optical characteristic is a chirp of a modulated laser, i.e., variations of the laser optical frequency with intentionally induced intensity modulation. In addition, the  
20 spectral phase measurements can be used to learn about dispersive properties of an optical fiber or other optical materials or components. Various optical analyzers have been developed to measure the amplitude and phase of optical signals.

[0003] Some optical analyzers for measuring the phase of optical signals require optical filtering, which typically involves using an optical grating. A  
25 concern with these optical analyzers is that the resolution of an optical grating is inherently limited and is directly dependent on the size of the grating. In addition, optical gratings are generally expensive, which increases the cost of the optical analyzers.

[0004] Other optical analyzers for measuring the phase of optical signals  
30 require complex signal processing calculations, such as multi-dimensional autocorrelation or cross-correlation calculations of optical fields. A concern with these optical analyzers is that sophisticated processors are needed to perform the complex calculations to measure the phase.

**[0005]** In view of the above-described concerns, there is a need for an optical analyzer and method for measuring spectral phase of optical signals that does not require optical filtering or complex calculations, such as multi-dimensional autocorrelation or cross-correlation calculations.

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## SUMMARY OF THE INVENTION

**[0006]** An optical analyzer and method for measuring optical properties of optical signals utilizes a heterodyne architecture to measure spectral amplitude and phase of a periodically modulated input optical signal, such as an optical signal from a periodically modulated distributed feedback (DFB) laser. The spectral amplitude and phase measurements are derived from a heterodyne signal, which is produced by combining and mixing the input optical signal and a local oscillator (LO) signal. The optical spectrum that is reconstructed from the heterodyne signal includes “inner” spectral peaks that contain phase information of the input optical signal. The inner spectral peaks may be produced by an optical or electrical mixing technique. The spectral phase of the input optical signal is recovered from the inner spectral peaks of the reconstructed optical spectrum.

**[0007]** The optical analyzer and method in accordance with the invention measures spectral amplitude and spectral phase. The spectral phase is determined from the spectral phase differences between adjacent spectral peaks of the input optical signal by using the inner spectral peaks. These measurements allow reconstruction of optical signals in the time domain by means of the Fourier transform. In addition, the direct measurement of the spectral phase allows dispersive properties of optical materials or components to be studied. The spectral phase differences are derived without the use of expensive optical filtering or complex calculations, such as multi-dimensional autocorrelation or cross-correlation calculations

**[0008]** Other aspects of the present invention will become apparent from the following detailed description, taken in conjunction with the accompanying drawings, illustrated by way of example of the principles of the invention.

## BRIEF DESCRIPTION OF THE DRAWINGS

[0009] Fig. 1 is a diagram of a heterodyne optical spectrum analyzer (HOSA) system in accordance with an embodiment of the invention.

5 [0010] Fig. 2 is an exemplary optical amplitude spectrum of a periodically modulated input optical signal.

[0011] Fig. 3 is an exemplary optical amplitude spectrum of a phase modulated local oscillator (LO) signal.

10 [0012] Fig. 4 is a simplified spectrum of a periodically modulated input optical signal.

[0013] Figs. 5A, 5B and 5C illustrate the mixing of the input optical signal and the phase modulated LO signal to produce a heterodyne signal with inner peaks.

15 [0014] Fig. 6 illustrates the concept of a triangle to solve for a phase difference between peaks of the input optical signal using the inner peaks of a heterodyne signal.

[0015] Fig. 7 illustrates a technique to measure the phase difference using the concept of Fig. 6 in accordance with a first embodiment of the invention.

20 [0016] Fig. 8A shows components of a modulation controller included in the HOSA system of Fig. 1 in accordance with the first embodiment of the invention.

[0017] Fig. 8B shows components of a processing unit included in the HOSA system of Fig. 1 in accordance with the first embodiment of the invention.

25 [0018] Figs. 9A and 9B illustrate a technique to measure the phase difference using the concept of Fig. 6 in accordance with a second embodiment of the invention.

[0019] Fig. 10A shows components of the modulation controller in accordance with the second embodiment of the invention.

[0020] Fig. 10B shows components of the processing unit in accordance with the second embodiment of the invention.

30 [0021] Fig. 11 is a diagram of a HOSA system in accordance with an alternative second embodiment of the invention.

[0022] Figs. 12A and 12B illustrate combining of input spectral peaks when using an electrical mixing technique in accordance with the alternative second embodiment of the invention.

5 [0023] Figs. 13A and 13B illustrate a technique to measure the phase difference using the concept of Fig. 6 in accordance with a third embodiment of the invention.

[0024] Fig. 14A shows components of the modulation controller in accordance with the third embodiment of the invention.

10 [0025] Fig. 14B shows components of the processing unit in accordance with the third embodiment of the invention.

[0026] Fig. 15 is a process flow diagram of a method for measuring optical properties of optical signals in accordance with an embodiment of the invention.

[0027] Fig. 16 is a process flow diagram of a method for measuring optical properties of optical signals in accordance with another embodiment of the  
15 invention.

#### DETAILED DESCRIPTION

[0028] With reference to Fig. 1, a heterodyne optical spectrum analyzer  
20 (HOSA) system 100 in accordance with an embodiment of the invention is shown. The HOSA system 100 operates to measure spectral amplitude and phase of an input optical signal, such as an optical signal from a periodically modulated distributed feedback (DFB) laser or other periodically modulated optical signal, using a heterodyne architecture. The measured spectral amplitude and phase  
25 allow the HOSA system 100 to compute time-varying optical characteristics of the input optical signal, such as amplitude modulation, phase modulation and chirp by means of the Fourier transform. Since the HOSA system 100 uses an optical heterodyne architecture, expensive optical filters and complex calculations, such as multi-dimensional autocorrelation and cross-correlations calculations of optical  
30 fields, which are required in some conventional systems, are not needed.

[0029] The HOSA system 100 takes advantage of the fact that the beat of two optical signals at different frequencies contains information regarding the phase difference of the two optical signals. It is well known that two acoustic

waves at different frequencies will pulse or “beat” at a frequency that is equal to the difference between the frequencies of the two acoustic waves. As an example, two acoustic waves at 30 Hertz and 34 Hertz will produce a beat frequency of 4 Hertz. The phase of the beat contains information about the phase difference  
 5 between the two acoustic waves. This beat phenomena also occurs for optical signals, and thus, the beat of two optical signals at different optical frequencies contain phase information of the two optical signals. Furthermore, the phase information is preserved in a heterodyne signal, which may be produced by combining and mixing the optical signals with a swept local oscillator (LO)  
 10 signal. Thus, the phase information contained in the beat of the heterodyne signal can be extracted to measure the phase difference of the optical signals. Using this fact, the HOSA system 100 measures the phase differences between spectral peaks of a periodically modulated input optical signal by analyzing a heterodyne signal produced by combining and mixing the input optical signal with a phase  
 15 modulated LO signal, as described in detail below.

**[0030]** As shown in Fig. 1, the HOSA system 100 includes an input 102, a modulated LO 104, an optical coupler 106, an optical receiver 108 and a processing unit 110. A periodically modulated input optical signal to be analyzed is received at the input 102. The input optical signal may be from an optical  
 20 device, such as a DFB laser that is modulated using a predefined frequency  $f_s$ , and, therefore, whose spectrum comprises multiple spectral peaks separated by  $f_s$ . Alternatively, the input optical signal may be from an optical fiber, which is carrying the periodically modulated optical signal. The modulation frequency  $f_s$  can be recovered from the input signal. Thus, it is assumed that the frequency of  
 25 modulation  $f_s$ , which is equal to the separation between the spectral peaks, is known. An exemplary optical amplitude spectrum 200 of the input optical signal is illustrated in Fig. 2. The optical amplitude spectrum 200 of the input optical signal includes peaks 202 separated by a frequency interval  $f_s$  equal to the modulating frequency. In Fig. 1, the input 102 is connected to the optical coupler  
 30 106 via an optical path 112 to transmit the received input optical signal to the optical coupler.

**[0031]** The modulated LO 104 is also connected to the optical coupler 106 via an optical path 114. The modulated LO signal generated by the modulated LO

104 has at least one sideband peak created by modulation. As shown in Fig. 1, the modulated LO 104 includes an unmodulated LO source 116, an optical modulator 118, and a modulation controller 120. The unmodulated LO source 116 is described herein as a swept LO source. However, the LO source 116 may be  
 5 another type of LO source that can change the frequency of the generated LO signal. Similarly, the optical modulator 118 is described herein as a phase modulator. However, the optical modulator 118 may be an intensity modulator. The LO source 116 generates a swept LO signal, which is, for example, phase modulated at a frequency of  $f_m$  by the optical modulator 118. The optical  
 10 modulator 118 is controlled by the modulation controller 120, which provides an electrical phase modulation signal having a frequency of  $f_m$ . The frequency of  $f_m$  is typically related to the separation between the optical spectral peaks of the input optical signal equal to  $f_s$ . As an example, the frequency of  $f_m$  may be half the frequency separation  $f_s$  between the optical spectral peaks of the input optical  
 15 signal, or a multiple integer of half the frequency separation  $f_s$ . The electrical signal at the frequency  $f_s$  may be supplied to the modulation controller 120 from an external source. As an example, the optical modulator 118 may be a Ti-indiffused LiNbO<sub>3</sub> phase modulator. An exemplary optical amplitude spectrum 300 of the phase modulated LO signal is illustrated in Fig. 3. The optical  
 20 amplitude spectrum 300 of the phase modulated LO signal includes a central peak 302 at the carrier frequency and multiple sideband peaks 304 separated from the carrier by multiples of  $f_m$ . Preferably, the modulation depth (defined below in the equation (1)) of the LO signal is small so that the central peak 302 and the nearest sideband peaks 304 are substantially larger than other (higher order) sideband  
 25 peaks (not shown).

**[0032]** Turning back to Fig. 1, the optical coupler 106 is further connected to the optical receiver 108 via an output optical path 122. The input optical signal and the phase modulated LO signal on the optical paths 112 and 114, respectively, are combined at the optical coupler 106 and transmitted to the optical receiver 108  
 30 on the output optical path 122. The combined optical signals are then detected by the optical receiver 108, which may be a square-law detector. The square-law detection leads to mixing of the combined optical signals and produces a heterodyne signal having a frequency in radio frequency (RF) range, which is

equal to the frequency difference of the combined signals. When a modulated LO having multiple sideband peaks is used, there may be more than one heterodyne signal. Dual heterodyne signals produce a beat whose AM demodulation provides a signal for a recovery of the spectral phase of the input optical signal. The optical receiver 108 also converts the detected optical signals into an electrical signal, i.e., current or voltage, which is processed by the processing unit 110 to measure desired optical characteristics of the periodically modulated input optical signal, such as amplitude and phase in the time domain. Although the components of the processing unit 110 are later shown and described as hardware components, the components of the processing unit 110 may be implemented in any combination of hardware, software and firmware.

**[0033]** The problem of measuring the phase of the periodically modulated input signal from the heterodyne signal is now presented in the following mathematical description. The electric field of the optical wave from the LO source 116 is sinusoidally modulated by the phase modulator 118 at the frequency  $f_m$ . The phase of the electrical modulation signal applied to the optical modulator 118,  $\psi$ , is controlled by the modulation controller 120 using, for example, a variable delay line. The phase modulated LO signal is described by:

$$e_0(t) = a_0 e^{j2\pi\nu_0 t + j\alpha \cos(2\pi f_m t + \psi) + j\varphi_0}, \quad (1)$$

where  $a_0$  is the amplitude of the electric field,  $\nu_0$  is the optical frequency of the LO signal, which is typically a swept frequency (i.e.,  $\nu_0 = \nu_0(t)$ ),  $\varphi_0$  is the phase term that also denotes the phase noise (i.e.,  $\varphi_0 = \varphi_0(t)$ ),  $\alpha$  is the modulation depth, and  $f_m$  is the frequency of the phase modulation. The sinusoidal modulation of the LO signal produces multiple spectral peaks whose amplitude and frequency can be determined from the series expansion:

$$e^{j\alpha \cos \xi} = \sum_{n=-\infty}^{\infty} j^n J_n(\alpha) e^{jn\xi}, \quad (2)$$

where  $J_n(\alpha)$  are Bessel functions.

**[0034]** In order to simplify the mathematical analysis, a spectrum under test 400 comprising just two peaks 402 and 404 is considered, as shown in Fig. 4, rather than the entire spectrum of the periodically modulated input signal, which includes many peaks, as illustrated in Fig. 2. The two peaks 402 and 404 of the

spectrum 400 have optical frequencies  $\nu_1$  and  $\nu_2$ , amplitudes  $a_1$  and  $a_2$ , and phases  $\varphi_1$  and  $\varphi_2$ , respectively. Both spectral peaks 402 and 404 are assumed to have the same phase noise, which is represented by  $\varphi_n$ . Thus, the electrical field of the optical input signal is described by:

$$e_s(t) = a_1 e^{j2\pi\nu_1 t + j\varphi_1 + j\varphi_n} + a_2 e^{j2\pi\nu_2 t + j\varphi_2 + j\varphi_n}. \quad (3)$$

The objective is to measure the phase difference  $\Delta\varphi = \varphi_2 - \varphi_1$  to determine the relative phase between the two peaks 402 and 404.

**[0035]** The combined optical waves of the periodically modulated input signal and the phase modulated LO signal can be represented by the sum:

$$e(t) = e_0(t) + e_s(t). \quad (4)$$

In the equation (4), the coupling coefficients and phase shift of the optical coupler 106 are omitted to provide a more lucid description. The intensity,  $i$ , at the optical receiver 108 is equal to  $e \cdot e^*$ , and thus, can be expressed as:

$$\begin{aligned} i = & a_0^2 + a_1^2 + a_2^2 + 2a_1a_2 \cos(2\pi\nu_1 t - 2\pi\nu_2 t + \varphi_1 - \varphi_2) + \\ & 2a_0a_1 \cos(2\pi\nu_1 t - 2\pi\nu_0 t - \alpha \cos(2\pi f_m t + \psi) - \varphi_0 + \varphi_n) + \\ & 2a_0a_2 \cos(2\pi\nu_2 t - 2\pi\nu_0 t - \alpha \cos(2\pi f_m t + \psi) - \varphi_0 + \varphi_n). \end{aligned} \quad (5)$$

The terms of interest in the equation (5) are the last two “mixing” terms, which are the terms that are derived from the mixing of the periodically modulated input signal and the phase modulated LO signal. The significance of these “mixing” terms is now described.

**[0036]** Figs. 5A, 5B and 5C illustrate the combining and mixing of the periodically modulated input signal and the phase modulated LO signal to produce the heterodyne signal and the related optical spectrum. In Fig. 5A, the original spectrum 400 of the input optical signal and the original spectrum 300 of the LO signal are shown. The phase modulated LO signal, which is typically a swept LO signal, is assumed to comprise only the carrier peak 302 and the two nearest sideband peaks 304. Each peak of the phase modulated LO signal creates its own image of the original input spectrum at the RF, as illustrated in Fig. 5B. The spectra from Fig. 5B are combined into a reconstructed RF spectrum 500 shown in Fig. 5C. The reconstructed spectrum 500 includes peaks 502 and 504, which correspond to the original peaks 402 and 404 of the input optical signal, respectively. Also, there are additional spectral peaks that are produced by the LO



sideband peaks on both sides of the peaks 502 and 504. If  $f_m \approx f_s / 2$ , then there are two nearly overlapping or overlapping peaks 506 in-between the peaks 502 and 504. These peaks 506 are formed by the LO sideband peaks 304 and the original peaks 402 and 404 of the input signal. The two nearly overlapping or  
5 overlapping peaks 506 will produce a single observable spectral peak, which will be referred to herein as an “inner peak”. Thus, the inner peak is a combination of the original peak 402 and 404 of the input signal. The inner peak is created by the dual heterodyne signal that produces a beat (amplitude modulated heterodyne signal). The phase of that beat is related to the phase difference  $\varphi_2 - \varphi_1$ .  
10 Therefore, the inner peak contains the phase difference,  $\varphi_2 - \varphi_1$ , of the original peaks 402 and 404 of the input signal.

[0037] The analysis with respect to the peaks 506 of the reconstructed spectrum 500 can be shown mathematically by expanding the mixing terms of the equation (5) into a series. The following series expansions can be used to expand  
15 the mixing terms.

$$\cos(\alpha \cos \xi) = J_0(\alpha) + 2 \sum_{j=1}^{\infty} (-1)^j J_{2j}(\alpha) \cdot \cos(2j\xi) \quad (6)$$

$$\cos(\alpha \cos \xi) = 2 \sum_{j=0}^{\infty} (-1)^j J_{2j+1}(\alpha) \cdot \cos((2j+1)\xi) . \quad (7)$$

[0038] The peaks 506 of the reconstructed spectrum 500 are identified by their frequency and their dependence on  $J_1(\alpha)$ :

$$p = 2a_0a_1J_1(\alpha)\sin(2\pi(\nu_1 - \nu_0 + f_m)t + \varphi_1 + \psi + \varphi_n - \varphi_0) + 2a_0a_2J_1(\alpha)\sin(2\pi(\nu_2 - \nu_1 - f_m)t + \varphi_2 - \psi + \varphi_n - \varphi_0), \quad (8)$$

where  $p$  denotes the amplitude of the inner peak. By choosing  $\nu_0 = (\nu_1 + \nu_2)/2$  (LO frequency in-between the measured peaks 402 and 404 of the input signal) and  $f_m = (\nu_2 - \nu_1)/2$  (the modulation frequency equal to a half of the frequency difference between the measured peaks of the input signal, i.e.,  $f_m = f_s / 2$ ), both  
25 peaks 506 are mixed to DC. This choice is not necessary, but it simplifies the mathematical form of the equation (8) to:

$$b = a_1 \sin(\varphi_1 + \psi + \varphi_n - \varphi_0) + a_2 \sin(\varphi_2 - \psi + \varphi_n - \varphi_0) , \quad (9)$$

where  $b = p/(2a_0J_1(\alpha))$  is a normalized amplitude of the inner peak. The equation (9) describes a simple trigonometric problem of a triangle, as shown in Fig. 6.

[0039] The amplitude of the inner peak can be computed by the addition of phasors having amplitudes  $a_1$  and  $a_2$ , as shown in Fig. 6. From the equation of a triangle, the following equation can be derived.

$$b^2 = a_1^2 + a_2^2 + 2a_1a_2 \cos(\Delta\phi - 2\psi). \quad (10)$$

It is important to note that the phase noise terms  $\phi_n$  and  $\phi_0$  are not in the equation (10). The term to be solved for in the equation (10) is  $\Delta\phi$ , which is the phase difference between the two peaks 402 and 404 of the input optical signal having a spectrum 400, i.e.,  $\Delta\phi = \phi_2 - \phi_1$ .

[0040] According to a delay property of the Fourier transform:

$$\begin{aligned} F(\omega) &\xrightarrow{F^{-1}} f(t) \\ e^{-j\omega t_0} F(\omega) &\xrightarrow{F^{-1}} f(t - t_0) \end{aligned}$$

Since the input optical signal is periodic in the time domain, the delay  $t_0$  and the corresponding time shift  $e^{j\omega t_0}$  is of no consequence. This simplifies methods for solving the equation (10). Instead of solving for  $\Delta\phi$ , the equation (10) can be solved for  $\Delta\phi + \psi_r$ , where  $\psi_r$  is a constant that is the same for all the measured inner peaks. Consequently, the angle  $\Delta\phi + \psi_r$  becomes the new  $\Delta\phi$  to be found.

[0041] In a first embodiment of the invention, the equation (10) is solved using discrete phase shifts  $\pm \Delta\psi$  of the LO modulation signal. In this embodiment, the LO signal is phase modulated such that  $f_m = f_s/2$ . First, the size of the inner peak is measured for an arbitrary reference phase  $\psi = \psi_r$  and, then, for two other phases  $\psi_{+,-} = \psi_r \pm \Delta\psi$ . Thus, three measurements are made to compute the phase difference  $\Delta\phi$ . The phase shifts  $\pm \Delta\psi$  are introduced by the modulation controller 120, as described further below. Since the actual value of  $\psi_r$  does not matter (a delay property of the Fourier transform), it is assumed that  $\psi_r = 0$ . From the equation (10), the following equations are obtained.

$$b_-^2 = a_1^2 + a_2^2 + 2a_1a_2 \cos(\Delta\phi + 2\Delta\psi) \quad (11a)$$

$$b_r^2 = a_1^2 + a_2^2 + 2a_1a_2 \cos(\Delta\varphi) \quad (11b)$$

$$b_+^2 = a_1^2 + a_2^2 + 2a_1a_2 \cos(\Delta\varphi - 2\Delta\psi) \quad (11c)$$

By subtracting the equation (11b) from the equations (11a) and (11c), and performing some trigonometric simplifications, the solution for  $\tan \Delta\varphi$  is:

$$5 \quad \tan \Delta\varphi = \frac{b_+^2 - b_-^2}{2b_r^2 - b_+^2 - b_-^2} \cdot \tan \Delta\psi \quad (12)$$

The term  $\tan \Delta\psi$  from the equation (12) is made equal to one by choosing  $\Delta\psi = \pi/4$ .

[0042] Graphically, the phase shift technique to solve for  $\Delta\varphi$  in the equation (10) can be illustrated as selecting points from a circle traced by the phasor  $\vec{a}_2$ , e.g., points 702, 704 and 706, as shown in Fig. 7. Using the measurements of  $b_+$ ,  $b_-$  and  $b_r$ , the phase difference  $\Delta\varphi$  can be solved for using the equation (12).

[0043] Turning now to Figs. 8A and 8B, the components of the modulation controller 120 and the processing unit 110 of the HOSA system in accordance with the first embodiment of the invention are shown. As described above and shown in Fig. 1, the modulation controller 120 is connected to the phase modulator 118 to provide a modulation signal so that the sideband peaks are added to the LO spectrum from the LO source 116. As shown in Fig. 8A, the modulation controller 120 includes a modulation signal generator 802 and an adjustable delay 804. The modulation signal generator 802 provides the modulation signal whose frequency is equal to  $f_s$ . The signal at the frequency  $f_s$  may be supplied to the modulation signal generator 802 from an external source. The adjustable delay 804 provides the needed phase shifts  $\pm \Delta\psi$ . The modulation signal is then transmitted to the phase modulator 118 to modulate the LO signal from the LO source 116. The modulated LO signal is combined with the input optical signal at the optical coupler 106.

[0044] As shown in Fig. 8B, the processing unit 108 includes a pre-processing section 806, which comprises a preamplifier 808, a power spectrum generator 810 and memory 812. The preamplifier 808 is connected to the optical receiver 108 to receive an electrical signal of photo-converted current generated

by the optical receiver, which represent the heterodyne signal produced by the mixing of the input optical signal and the phase modulated LO signal. The preamplifier 808 operates to amplify the electrical signal from the optical receiver 108. The amplified electrical signal is then squared by the power spectrum generator 810 to produce a power spectrum, which is a reconstructed optical spectrum of the input optical signal having spectral peaks that correspond to the original spectral peaks of the input optical signal, as well as inner additional spectral peaks. The power spectrum is stored in the memory 812 for subsequent processing to compute the spectral amplitude and the spectral phase of the input optical signal. Depending on the number of measurements to be made using different phase shifts (three in this embodiment), a corresponding number of power spectra are generated and stored.

[0045] The processing unit 110 further includes a peak identifier 814, an amplitude computer 816 and a phase computer 818. For each power spectrum, the peak identifier 814 determines the peaks of the power spectrum, including the original spectral peaks and the inner spectral peaks. The peaks that correspond to the original spectral peaks of the input optical signal are used to compute the spectral amplitude of the input signal by the amplitude computer 816. Since multiple measurements are not needed to compute the spectral amplitude, only the respective peaks from one of the stored power spectrum may be used.

Alternatively, multiple power spectra may be used to find an average measurement of the amplitude. The corresponding inner peaks of the three power spectra are used to compute the phase difference,  $\Delta\phi$ , by the phase computer 818 using the equation (12). The spectral phase is computed from the phase difference by summing consecutive phase differences, i.e.,  $\phi_i = \sum \Delta\phi_i$ . The computed spectral amplitude and phase of the input optical signal can then be further processed using the Fourier transform to compute the amplitude and phase in the time domain. The chirp of the input optical signal can be found from the derivative of phase in the time domain.

[0046] In a second embodiment of the invention, the phase difference  $\Delta\phi$  is measured by comparing the phase of the oscillatory behavior of the inner peaks to the phase of the reference signal at the electrical frequency  $f_r = \delta f$ . A graphical

explanation of this phase difference measurement is now described. Assume that the peaks 506 in Fig. 5C are not at the same frequency but that there is a frequency difference between the peaks equal to  $\delta f$ . This means that the phase modulation frequency is not equal to  $f_s/2$  but to  $f_s/2 + \delta f/2$  or to

5  $f_s/2 - \delta f/2$ . Then, using the concept of Fig. 6, the phasor  $\vec{a}_2$  will rotate around the phasor  $\vec{a}_1$ , as shown in Fig. 9A. Therefore, the amplitude of the resulting inner peak, which is proportional to  $b$ , will oscillate at the frequency  $\delta f$ , as shown in Fig. 9B. The oscillatory behavior of  $b(t)$  can then be defined by the following equation:

$$10 \quad b(t)^2 = A + B \cos(2\pi\delta f t + \Delta\psi), \quad (13)$$

where  $A$  and  $B$  are some constants. The phase of the oscillating  $b(t)^2$  from the equation (13), as compared to the phase of the reference signal  $r(t) = C \cos(2\pi\delta f t)$ , is a measure of the phase difference  $\Delta\phi$ .

[0047] Turning now to Figs. 10A and 10B, the components of the modulation controller 120 and the processing unit 110 of the HOSA system 100 in accordance with the second embodiment of the invention is shown. In this embodiment, the modulation controller 120 includes only a modulation signal generator 1002 and does not include an adjustable delay, as shown in Fig. 10A. The modulation signal generator 1002 provides the modulation signal that is offset from the frequency  $f_s/2$  by  $\delta f/2$ , i.e.,  $f_m = f_s/2 + \delta f/2$  or  $f_m = f_s/2 - \delta f/2$ . The electrical signal at the frequency  $f_s$  may be supplied to the modulation signal generator 1002 from an external source. The modulation signal is transmitted to the phase modulator 118 to modulate the LO signal from the LO source 116. The modulated LO is combined with the input optical signal at the optical coupler 106.

25 [0048] As shown in Fig. 10B, the processing unit 110 includes a preamplifier 1004, an AM demodulator 1006 and a phase sensitive detector 1008. The preamplifier 1004 amplifies the electrical signal generated by the optical receiver 108, i.e., the heterodyne signal that can be used to produce a reconstructed optical spectrum of the input optical signal. The amplified electrical signal is then demodulated by the AM demodulator 1006. The AM demodulator 1006 recovers the amplitude of the spectral peaks of the input optical signal.

Then, the demodulated signal is transmitted to the phase sensitive detector 1008, where the demodulated signal containing the beat between the two heterodyne signals of the inner spectral peaks (their amplitude modulation) is compared with a reference signal having a frequency of  $f_r = \delta f$  to recover the spectral phase difference of the input optical signal for all pairs of the adjacent sideband peaks. The reference frequency  $f_r$  can be produced from  $f_m$  and  $f_s$  by mixing the two frequencies. The measured spectral amplitude and the spectral phase can then be further processed to compute the amplitude and phase of the input optical signal in the time domain by means of the Fourier transform. The chirp of the input optical signal can be found from the derivative of the computed phase in the time domain.

[0049] In Fig. 11, a HOSA system 1100 in accordance with an alternative second embodiment is shown. The HOSA system 110 uses the rotating phasor concept of Figs. 9A and 9B to derive spectral amplitude and phase of an input optical signal. However, the HOSA system 1100 is configured to use an electrical mixing technique, rather than an optical mixing technique, so that an unmodulated LO signal can be used. Thus, the phase modulator 118 and the modulation controller 120 are not included in the HOSA system 1100. In this alternative embodiment, the bandwidth of the optical receiver 108 must be wider than  $f_s$ .

[0050] As shown in Fig. 11, the HOSA system 1100 includes a processing unit 1102, which comprises a preamplifier 1104, a mixer 1106, an AM demodulator 1108, a phase sensitive detector 1110 and an amplitude calculator 1112. The preamplifier 1104 amplifies the electrical signal generated by the optical receiver 108, i.e., the heterodyne signal. The amplified electrical signal is mixed at the mixer 1106 with an electrical signal having a frequency  $f_d$ , which is equal to  $f_s/2 - \delta f$  or  $f_s/2 + \delta f$ . As illustrated in Figs. 12A and 12B, this results in a simultaneous observation of the adjacent spectral peaks 1202 and 1204 of the input signal when the LO signal 1206 is approximately in-between the peaks of the optical input signal. Fig. 12A shows the optical LO signal 1206 being in-between the adjacent optical spectral peaks 1202 and 1204 of the optical input signal. The optical LO signal 1206 combines with each of the adjacent optical spectral peaks 1202 and 1204 to produce the electrical spectral peaks 1208 and 1210 in the RF (shown in Fig. 12B) that correspond to the original adjacent

spectral peaks. Fig. 12B shows the RF spectral peaks 1208 and 1210 and the electrical signal 1212 having frequencies  $\pm f_d$  that are electrically mixed to produce dual heterodyne signal being amplitude modulated at the frequency  $\delta f$ . The dual heterodyne signal forms the desired “inner” spectral peak that combines the adjacent spectral peaks of the original spectrum of the input optical signal.

[0051] The phase of the AM of the dual heterodyne signal contains the information about the phase difference. In addition, the dual heterodyne signal is used to reconstruct spectral peaks that are related to the original spectral peaks of the input optical signal. Namely, each reconstructed peak is a combination of two adjacent peaks of the original spectrum. Thus, each reconstructed spectral peak is an inner peak that contains phase difference information of the adjacent original spectral peaks and power of the adjacent original peaks, i.e., on average

$b_i^2 = a_i^2 + a_{i+1}^2$  as in the equation (10). In other words, the oscillatory behavior of  $b_i$  described by the equation (10) provides phase difference information, while its

average amplitude can be used to find the original amplitudes  $a_i$  (original spectral amplitude). The mixed electrical signal is then demodulated by the AM demodulator 1108. The AM demodulator 1108 recovers the amplitude of the inner spectral peaks that is used to find the amplitude of the original spectral peaks of the input optical signal from the equation  $b_i^2 = a_i^2 + a_{i+1}^2$  by the amplitude

calculator 1112, which in effect reconstructs the original spectral peaks of the optical input signal from the inner spectral peaks. The demodulated signal is also transmitted to the phase sensitive detector 1110, where the demodulated signal is compared with a reference signal having a frequency  $f_r = f_s - 2f_d = \delta f$  to

recover the spectral phase of the input signal. The reference signal at the

frequency  $f_r$  can be constructed from  $f_s$  and  $2f_d$  by mixing the two frequencies.

The measured spectral amplitude and the spectral phase can then be further processed to compute the amplitude and phase of the input optical signal in the time domain by means of the Fourier transform. The chirp of the input optical signal can be found from the derivative of the computed phase in the time domain.

[0052] In a third embodiment of the invention, the phase difference  $\Delta\phi$  is measured by monitoring an oscillatory behavior of the inner peaks caused by

phase modulation of the electrical modulation signal applied to the optical modulator 118. An explanation of the phase shift measurement in accordance with the third embodiment is now described. Assume that the modulation controller 120 from Fig. 1 provides a signal at the electrical frequency  $f_s$  that is  
 5 phase modulate by a phase term  $\psi(t) = \delta\psi \cos(\omega_\psi t + \Theta)$ , where  $\omega_\psi = 2\pi f_\psi$ .  
 Then, using the concept of Fig. 6, the phasor  $\vec{a}_2$  will move back and forth such that the angle between the phasors  $\vec{a}_1$  and  $\vec{a}_2$  will continuously vary in a sinusoidal manner, as shown in Fig. 13A. Consequently, the amplitude of the phasor  $\vec{b}$  will change periodically, as shown in Fig. 13B, according to the  
 10 following function.

$$b(t)^2 = A + B \cos(\Delta\phi + \delta\psi \cos(\omega_\psi t + \Theta)), \quad (14)$$

where A and B are constants. Thus,  $b(t)$  contains harmonics of the electrical phase modulation frequency  $f_\psi$ . An appropriate series expansion of the function (14) gives:

$$\begin{aligned} b(t)^2 = & BJ_0(\delta\psi) \cos(\Delta\phi) \\ & - 2BJ_2(\delta\psi) \cos(2(\omega_\psi + \Theta)) \cos(\Delta\phi) \\ & + 2BJ_4(\delta\psi) \cos(4(\omega_\psi + \Theta)) \cos(\Delta\phi) \\ & \vdots \\ 15 \quad & - 2BJ_1(\delta\psi) \cos(\omega_\psi + \Theta) \sin(\Delta\phi) \\ & + 2BJ_3(\delta\psi) \cos(3(\omega_\psi + \Theta)) \sin(\Delta\phi) \\ & - 2BJ_5(\delta\psi) \cos(5(\omega_\psi + \Theta)) \sin(\Delta\phi) \\ & \vdots \end{aligned} \quad (15)$$

The relative amplitudes of the odd and even harmonics are proportional to  $\sin(\Delta\phi)$  and  $\cos(\Delta\phi)$ , respectively. Thus, measurements of one odd harmonic and one even harmonic allow the phase difference  $\Delta\phi$  to be found. For example, the phase difference  $\Delta\phi$  may be computed by using the amplitude of the first  
 20 harmonic,  $h_1 = 2BJ_1(\delta\psi) \sin(\Delta\phi)$ , and the amplitude of the second harmonic,  $h_2 = 2BJ_2(\delta\psi) \cos(\Delta\phi)$ , in accordance with the following equation:

$$\tan \Delta\phi = \frac{h_1 J_2(\delta\psi)}{h_2 J_1(\delta\psi)} \quad (16)$$



[0053] Turning now to Figs. 14A and 14B, the components of the modulation controller 120 and the processing unit 110 of the HOSA system 100 in accordance with the third embodiment of the invention is shown. In this embodiment, the modulation controller 120 includes a modulation signal generator 1402 and an electrical frequency modulator 1404, as shown in Fig. 14A. The modulation signal generator 1402 provides a modulation signal at the electrical frequency  $f_s$ . The electrical frequency  $f_s$  may be supplied from an external source. The modulation signal at the electrical frequency  $f_s$  is then phase modulated by the electrical frequency modulator 1404 at the electrical frequency  $f_\psi$  and supplied to the optical modulator 118.

[0054] As shown in Fig. 14B, the processing unit 110 includes a preamplifier 1406, an AM demodulator 1408, a power converter 1410, phase sensitive detectors 1412 and 1414 and a processor 1416. The preamplifier 1406 amplifies the electrical signal generated by the optical receiver 108, i.e., the heterodyne signal that can be used to produce a reconstructed optical spectrum of the input optical signal. The amplified electrical signal is then demodulated by the AM demodulator 1408. The AM demodulator 1408 recovers the amplitude of the spectral peaks of the input optical signal. The demodulated signal is then squared by the power converter 1410 and transmitted to the phase sensitive detectors 1412 and 1414. The amplitudes of odd harmonics and even harmonics of the electrical phase modulation frequency  $f_\psi$  are measured at the phase sensitive detectors 1412 and 1414. As an example, the phase sensitive detector 1412 may be configured to measure the amplitude of the first harmonics,  $h_1$ , using a reference signal having an electrical frequency of  $f_\psi$ , while the phase sensitive detector 1414 is configured to measure the amplitude of the second harmonic,  $h_2$ , using another reference signal having an electrical frequency of  $2f_\psi$ . These measured amplitudes are transmitted to the processor 1416 where the phase difference  $\Delta\phi$  is computed using the equation (16) to determine the spectral phase of the optical input signal. The recovered spectral amplitude and the spectral phase difference can then be further processed to compute the amplitude and phase of the input optical signal in the time domain by means of the Fourier transform. The chirp of the input optical signal can be found from the derivative of the computed phase in the time domain.

[0055] A method for analyzing optical properties of optical signals, such as amplitude and phase in the time domain, in accordance with an embodiment of the invention is described with reference to a flow diagram of Fig. 15. At block 1502, an input optical signal is received. The input optical signal is a periodically  
5 modulated signal. Thus, the input optical signal includes peaks at frequencies separated by fixed frequency intervals. At block 1504, a modulated LO signal having a central peak and sideband peaks is provided. Next, at block 1506, the input optical signal and the modulated LO signal are combined. At block 1508, the combined signals are mixed to construct spectral peaks that are combinations  
10 of the spectral peaks of the input optical signal. Next, at block 1510, the spectral phase differences between the spectral peaks of the input optical signal are derived using the constructed spectral peaks.

[0056] A method for analyzing optical properties of optical signals, such as amplitude and phase in the time domain, in accordance with another embodiment  
15 of the invention is described with reference to a flow diagram of Fig. 16. At block 1602, an input optical signal is received. The input optical signal is a periodically modulated signal. Thus, the input optical signal includes peaks at frequencies separated by fixed frequency separations. At block 1604, an unmodulated LO signal is provided. Next, at block 1606, the input optical signal and the LO signal  
20 are combined. At block 1608, the combined signals are mixed to produce a heterodyne signal. Next, at block 1610, the heterodyne signal is electrically mixed with an electrical signal to produce a mixed electrical signal having spectral peaks that are combinations of the spectral peaks of the input optical signal. At  
block 1612, the mixed electrical signal is compared with a reference signal to  
25 measure the spectral phase differences between the spectral peaks of the input optical signal.

[0057] Although specific embodiments of the invention have been described and illustrated, the invention is not to be limited to the specific forms or arrangements of parts so described and illustrated. The scope of the invention is  
30 to be defined by the claims appended hereto and their equivalents.